

GRAY-HERVELLA CLASSES ON PRODUCT TWISTOR SPACES

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ABSTRACT. Motivated by generalized geometry (in the sense of Hitchin), the product bundle $\mathcal{Z} \times_M \mathcal{Z}$ of the twistor space $\mathcal{Z}$ of a Riemannian manifold (M, g) is considered. The product twistor space admits a natural family of Riemannian metrics and four compatible almost complex structures, analogs of the Atiyah-Hitchin-Singer and Eells-Salamon almost complex structures on the twistor space. The Gray-Hervella classes of these almost Hermitian structures are determined in the case when the dimension of the base manifold M is four.

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1. INTRODUCTION

The Gray-Hervella classification [11] of almost Hermitian structures is widely recognized as a useful tool in studying these structures. In the present paper, the possible Gray-Hervella classes are determined on the product bundle $\mathcal{Z} \times_M \mathcal{Z} \rightarrow M$ (with fibre the product of fibres at the same point) of the twistor space $\mathcal{Z}$ of a Riemannian manifold (M, g) . Recall that the latter space is the bundle over M whose fibre at a point $p \in M$ consists of g -orthogonal complex structures on the tangent space $T_p M$. The motivation for considering the product twistor space $\mathcal{Z} \times_M \mathcal{Z}$ comes from the generalized complex geometry in the sense of N. Hitchin [13, 12]. Given a generalized Riemannian metric on a smooth manifold M , one can consider the bundle $\mathcal{G}$ over M whose fibre at a point p is the space of generalized complex structures on $T_p M$ compatible with the generalized metric. This bundle, called the generalized twistor space of M , is studied in [6]. The generalized metric yields a Riemannian metric g on M and a g -metric connection D with totally skew-symmetric torsion. If $\mathcal{Z}$ is the usual twistor space of the Riemannian manifold (M, g) , then the generalized twistor space $\mathcal{G}$ is diffeomorphic to the product twistor space $\mathcal{Z} \times_M \mathcal{Z}$. In [6], using this observation and the connection D , four generalized almost complex structures are defined on the generalized twistor space. One of them is an analog of the Atiyah-Hitchin-Singer [1] almost complex structure, the others of the Eells-Salamon [8] one on the usual twistor space. These generalized almost complex structures are not induced by usual almost complex structures, although the generalized twistor space admits four almost complex structures also defined by means of the diffeomorphism $\mathcal{G} \cong \mathcal{Z} \times_M \mathcal{Z}$ and the connection D . The integrability conditions for the four generalized almost complex structures are found in [6], the integrability conditions for the four usual almost complex structures are established in [7] in the case when the dimension of M is four, the most interesting

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case. These almost complex structures are compatible with a natural family of Riemannian metrics on $Z \times_M Z$. In this paper, the Gray-Hervella classes of the corresponding almost Hermitian structures are found under the natural assumption that the torsion of the connection D vanishes, i.e. D is the Levi-Civita connection of g (this is one of the integrability conditions for the analog of the Atiyah-Hitchin-Singer structure on $\mathcal{G}$).

2. BASICS OF TWISTOR SPACES

2.1. The manifold of compatible linear complex structures. Let T be a $2m$ -dimensional real vector space endowed with an Euclidean metric g . Recall that a complex structure J on T is called compatible with g if it is g -orthogonal: $g(JX, JY) = g(X, Y)$. The set $Z = Z(T, g)$ of all compatible complex structures on T has a natural smooth structure of an embedded submanifold of the vector space $\mathfrak{so}(g)$ of skew-symmetric endomorphisms of (T, g) .

The group of orthogonal transformations of (T, g) acts smoothly and transitively on the set Z by conjugation. The isotropy subgroup at a fixed $J \in Z$ consists of the orthogonal transformations commuting with J . Therefore Z can be identified with the homogeneous space $O(2m)/U(m)$. In particular, $\dim Z = m^2 - m$. Note also that Z has two connected components. If we fix an orientation on T , these components consist of all complex structures on T compatible with the metric g and inducing $\pm$ the orientation of T ; each of them has the homogeneous representation $SO(2m)/U(m)$.

The tangent space at a point $J \in Z$ is $T_J Z = \{V \in \mathfrak{so}(g) : JV + VJ = 0\}$. The smooth manifold Z admits an almost complex structure $\mathcal{K}$ defined by $\mathcal{K}V = J \circ V$. The almost complex structure $\mathcal{K}$ is compatible with the restriction G to Z of the standard metric $-\frac{1}{2}\text{Trace } AB$ of $\mathfrak{so}(g)$.

Given an orthonormal basis $e_1, \dots, e_{2m}$ of T , denote by $S_{a,b}$, $a, b = 1, \dots, 2m$, the skew-symmetric endomorphisms of T defined by

$$S_{a,b}e_c = \delta_{ac}e_b - \delta_{bc}e_a, \quad c = 1, \dots, 2m.$$

Then $\{S_{ab} : a < b\}$ is a G -orthonormal basis of $\mathfrak{so}(g)$.

Let $J \in Z$ and let $e_1, \dots, e_{2m}$ be an orthonormal basis of T such that $Je_{2i-1} = e_{2i}$, $i = 1, \dots, m$. Set

$$A_{r,s} = \frac{1}{\sqrt{2}}(S_{2r-1,2s-1} - S_{2r,2s}), \quad B_{r,s} = \frac{1}{\sqrt{2}}(S_{2r-1,2s} + S_{2r,2s-1}),$$

$$r = 1, \dots, m-1, \quad s = r+1, \dots, m.$$

Then $\{A_{r,s}, B_{r,s}\}$ is a G -orthonormal basis of $T_J Z$ with $B_{r,s} = \mathcal{K}A_{r,s}$.

Denote by D the Levi-Civita connection of the metric G on Z . Let X, Y be vector fields on Z considered as $\mathfrak{so}(g)$ -valued functions on $\mathfrak{so}(g)$. By the Koszul formula for the Levi-Civita connection, for every $J \in Z$,

$$(1) \quad (D_X Y)_J = \frac{1}{2}(Y'(J)(X_J) + J \circ Y'(J)(X_J) \circ J)$$

where $Y'(J) \in \text{Hom}(\mathfrak{so}(g), \mathfrak{so}(g))$ is the derivative of the function $Y : \mathfrak{so}(g) \rightarrow \mathfrak{so}(g)$ at the point J . The latter formula easily implies that $D\mathcal{K} = 0$, so $(G, \mathcal{K})$ is a Kähler structure on Z . Note also that the metric G is Einstein with scalar curvature $\frac{m}{2}(m-1)^2$ (see, for example, [4]).

In the case when T is of dimension four, each of the two connected components Z can be identified with a 2-sphere. It is often convenient to describe this identification in terms of the space $\Lambda^2 T$.

The metric g of T induces a metric on $\Lambda^2 T$ given by

$$(2) \quad g(x_1 \wedge x_2, x_3 \wedge x_4) = g(x_1, x_3)g(x_2, x_4) - g(x_1, x_4)g(x_2, x_3).$$

Consider the isomorphism $\mathfrak{so}(g) \cong \Lambda^2 T$ sending $\varphi \in \mathfrak{so}(g)$ to the 2-vector $\varphi^\wedge$ for which

$$g(\varphi^\wedge, x \wedge y) = g(\varphi x, y), \quad x, y \in \mathbb{R}^4.$$

This isomorphism is an isometry with respect to the metric G on $\mathfrak{so}(g)$ and the metric g on $\Lambda^2 \mathbb{R}^4$.

Fix an orientation on T and denote by $Z_\pm$ the set of complex structures on T compatible with the metric g and inducing $\pm$ the orientation of T .

If $\dim T = 4$, the Hodge star operator defines an endomorphism $*$ of $\Lambda^2 T$ with $*^2 = Id$. Hence we have the orthogonal decomposition

$$\Lambda^2 T = \Lambda_-^2 T \oplus \Lambda_+^2 T,$$

where $\Lambda_\pm^2 T$ are the subspaces of $\Lambda^2 T$ corresponding to the (± 1) -eigenvalues of the operator $*$. Let (e_1, e_2, e_3, e_4) be an oriented orthonormal basis of $\mathbb{R}^4$. Set

$$(3) \quad s_1^\pm = \frac{1}{\sqrt{2}}(e_1 \wedge e_2 \pm e_3 \wedge e_4), \quad s_2^\pm = \frac{1}{\sqrt{2}}(e_1 \wedge e_3 \pm e_4 \wedge e_2), \quad s_3^\pm = \frac{1}{\sqrt{2}}(e_1 \wedge e_4 \pm e_2 \wedge e_3).$$

Then $(s_1^\pm, s_2^\pm, s_3^\pm)$ is an orthonormal basis of $\Lambda_\pm^2 \mathbb{R}^4$.

The isomorphism $\varphi \rightarrow \varphi^\wedge$ identifies $Z_\pm$ with the sphere of radius $\sqrt{2}$ in the Euclidean vector space $(\Lambda_\pm^2 T, g)$. Under this isomorphism, if $J \in Z_\pm$, the tangent space $T_J Z = T_J Z_\pm$ is identified with the orthogonal complement $(\mathbb{R}J)^\perp$ of the space $\mathbb{R}J$ in $\Lambda_\pm^2 \mathbb{R}^4$.

The orientation of $\Lambda_\pm^2 T$ defined by the basis (3) does not depend on the choice of the basis $(e_1, \dots, e_4)$ ([5]). Consider the 3-dimensional Euclidean space $(\Lambda_\pm^2 T, g)$ with this orientation and denote by $\times$ the usual vector-cross product on it. For $\sigma, \tau \in \Lambda_\pm^2 T$, denote by K_σ and K_τ the endomorphisms of T corresponding to σ and τ under the isomorphism $\Lambda^2 T \cong \mathfrak{so}(g)$. Then $\sigma \times \tau$ corresponds to $\pm \frac{1}{\sqrt{2}}[K_\sigma, K_\tau]$. Thus, if $J \in Z_\pm$ and $V \in T_J Z = T_J Z_\pm$,

$$(\mathcal{K}V)^\wedge = \pm \frac{1}{\sqrt{2}}(J^\wedge \times V^\wedge).$$

2.2. The twistor space of a Riemannian manifold. Let (M, g) be a Riemannian manifold, $\dim M = 2m$. Denote by $\mathcal{Z} = \mathcal{Z}(TM, g)$ the bundle over M whose fibre at a point $p \in M$ is $Z(T_p M, g_p)$, the space of complex structures on $T_p M$ compatible with the metric g_p (the usual twistor space of (M, g)). Consider $\mathcal{Z}$ as a subbundle of the bundle $\mathcal{A} = \mathcal{A}(TM, g)$ of g -skew-symmetric endomorphisms of TM . Denote the projection of the bundle $\mathcal{A}$ onto M and its restriction to $\mathcal{Z}$ by π . Then the vertical space $\mathcal{V}_J$ of $\mathcal{Z}$ at a point $J \in \mathcal{Z}$ consists of skew-symmetric endomorphisms V of $(T_{\pi(J)}, g)$ anti-commuting with J . Let ∇ be the Levi-Civita connection of (M, g) . The induced connection on the bundle $\mathcal{A}$ will also be denoted by ∇ . Let $\mathcal{H}_J$ be the horizontal subspace of $T_J \mathcal{A}$ with respect to ∇ . Take an orthonormal basis $e_1, \dots, e_{2m}$ of $T_p M$, $p = \pi(J)$, such that $Je_{2k-1} = e_{2k}$, $Je_{2k} = -e_{2k-1}$, $k = 1, \dots, m$. Extend this basis to a frame of vector fields $E_1, \dots, E_{2m}$ such that $\nabla E_i|_p = 0$ and

define a section K of $\mathcal{A}$ by $KE_{2k-1} = E_{2k}$, $KE_{2k} = -E_{2k-1}$. Obviously K takes values in $\mathcal{Z}$ and $\nabla K|_p = 0$. Hence the horizontal space $\mathcal{H}_J = K_*(T_p M)$ is tangent to $\mathcal{Z}$. Thus we have the decomposition $T\mathcal{Z} = \mathcal{V} \oplus \mathcal{H}$. Using this decomposition, we can define a 1-parameter family h_t , $t > 0$, of Riemannian metrics on $\mathcal{Z}$ by

$$h_t(X^h + V, Y^h + W) = g(X, Y) + tG(V, W)$$

where X^h, Y^h are the horizontal lifts of $X, Y \in TM$, V, W are vertical vectors, and G is (the restriction) of the standard metric $-\frac{1}{2}\text{Trace}_g AB$ of $\mathcal{A}$. By the Vilms theorem [16] (or [2, Theorem 9.59]) the projection map $\pi : (\mathcal{Z}, h_t) \rightarrow (M, g)$ is a Riemannian submersion with totally geodesic fibres (this can also be proved by a direct computation). Recall also that the metrics h_t are compatible with the Atiyah-Hitchin-Singer [1] and Eeels-Salamon [8] almost complex structures $\mathcal{J}_1$ and $\mathcal{J}_2$ defined by

$$\mathcal{J}_1 X_J^h = \mathcal{J}_2 X_J^h = (JX)_J^h, \quad \mathcal{J}_1 V = -\mathcal{J}_2 V = J \circ V, \quad X \in T_{\pi(J)} M, \quad V \in \mathcal{V}_J.$$

2.3. The product twistor space. Now, let $\mathcal{G}$ be the product bundle $\mathcal{Z} \times_M \mathcal{Z}$. We shall consider $\mathcal{G}$ as a subbundle of the vector bundle $\pi : \mathcal{A} \oplus \mathcal{A} \rightarrow M$ endowed with the connection induced by the Levi-Civita connection of (M, g) , both denoted by ∇ . The vertical space of $\mathcal{G}$ at a point $J = (J_1, J_2)$ consists of the pairs (V_1, V_2) of skew-symmetric endomorphisms V_1, V_2 of $(T_{\pi(J)} M, g)$ anti-commuting with J_1, J_2 , respectively. The horizontal spaces of $\mathcal{A} \oplus \mathcal{A}$ are tangent to $\mathcal{G}$ and we have the decomposition $T\mathcal{G} = \mathcal{H} \oplus \mathcal{V}$.

For $t_1 > 0$, $t_2 > 0$, set $\mathbf{t} = (t_1, t_2)$ and define a Riemannian metric on the manifold $\mathcal{G}$ by

$$H_{\mathbf{t}}(X^h + V, Y^h + W) = g(X, Y) + t_1 G(V_1, W_1) + t_2 G(V_2, W_2)$$

where $X, Y \in TM$ and $V = (V_1, V_2)$, $W = (W_1, W_2)$ are vertical vectors. Then the projection map $\pi : (\mathcal{G}, H_{\mathbf{t}}) \rightarrow (M, g)$ is a Riemannian submersion.

Any fibre $\mathcal{V}_J$ of $\mathcal{G}$ possesses four complex structures $\mathcal{K}_n$ for which $(G_{\mathbf{t}}|_{\mathcal{V}_J}, \mathcal{K}_n)$ are Kähler structures on the fibre. These are defined by

$$\mathcal{K}_1(V_1, V_2) = (J_1 \circ V_1, J_2 \circ V_2), \quad \mathcal{K}_2(V_1, V_2) = (J_1 \circ V_1, -J_2 \circ V_2)$$

$$\mathcal{K}_3 = -\mathcal{K}_2, \quad \mathcal{K}_4 = -\mathcal{K}_1.$$

Correspondingly, we can define four almost complex structures $\mathcal{J}^n$ on $\mathcal{G}$ setting for $J = (J_1, J_2) \in \mathcal{G}$, $X \in T_{\pi(J)} M$, and $V \in \mathcal{V}_J$

$$\mathcal{J}^n X_J^h = (J_1 X)_J^h, \quad \mathcal{J}^n V = \mathcal{K}_n V, \quad n = 1, \dots, 4.$$

Clearly, these almost complex structures are compatible with the metrics $H_{\mathbf{t}}$.

3. TECHNICAL LEMMAS

Let $(\mathcal{U}, x_1, \dots, x_{2m})$ be a local coordinate system of M and $\{E'_1, \dots, E'_{2m}\}, \{E''_1, \dots, E''_{2m}\}$ orthonormal frames of TM on $\mathcal{U}$, respectively. Define sections S'_{ij}, S''_{ij} , $1 \leq i, j \leq 2m$, of the bundle $\mathcal{A}$ by the formulas

$$(4) \quad S'_{ij} E'_k = \delta_{ik} E'_j - \delta_{kj} E'_i, \quad S''_{ij} E''_k = \delta_{ik} E''_j - \delta_{kj} E''_i,$$

For $a = (a', a'') \in \mathcal{A} \oplus \mathcal{A}$, set

$$(5) \quad \tilde{x}_i(a) = x_i \circ \pi(a), \quad y'_{kl}(a) = G(a', S'_{kl} \circ \pi(a)), \quad y''_{kl}(a) = G(a'', S''_{kl} \circ \pi(a)).$$

Then $((\tilde{x}_i), (y'_{jk}), (y''_{rs}))$, $1 \leq i \leq 2m$, $1 \leq j < k \leq 2m$, $1 \leq r < s \leq 2m$, is a local coordinate system on the total space of the bundle $\mathcal{A} \oplus \mathcal{A}$.

If

$$(6) \quad V = \sum_{1 \leq j < k \leq 2m} [v'_{jk} \frac{\partial}{\partial y'_{jk}}(J) + v''_{jk} \frac{\partial}{\partial y''_{jk}}(J)].$$

is a vertical vector of $\mathcal{G}$ at a point J , the vertical vector $\mathcal{J}^n V$ is given by

$$(7) \quad \mathcal{J}^n V = (-1)^{n+1} \sum_{j < k} \sum_{s=1}^{2m} [\pm v'_{js} y'_{sk} \frac{\partial}{\partial y'_{jk}} + v''_{js} y''_{sk} \frac{\partial}{\partial y''_{jk}}]$$

where the plus sign corresponds to $n = 1, 4$ and the minus sign to $n = 2, 3$.

Note also that, for every $X \in TM$, we have

$$(8) \quad \mathcal{J}^n X^h = \sum_{i,j=1}^{2m} (g(X, E'_i) \circ \pi) y'_{ij} E'_j{}^h.$$

It is convenient to set $v'_{ij} = -v'_{ji}$, $v''_{ij} = -v''_{ji}$ and $y'_{ij} = -y'_{ji}$, $y''_{ij} = -y''_{ji}$ for $i \geq j$, $1 \leq i, j \leq 2m$. Then the endomorphism V of $T_p M \oplus T_p M$, $p = \pi(J)$, is determined by

$$V E'_i = \sum_{j=1}^{2m} v'_{ij} E'_j, \quad V E''_i = \sum_{j=1}^{2m} v''_{ij} E''_j.$$

For a vector field

$$X = \sum_{i=1}^{2m} X^i \frac{\partial}{\partial x_i}$$

on $\mathcal{U}$, the horizontal lift X^h on $\pi^{-1}(\mathcal{U})$ is given by

$$(9) \quad X^h = \sum_{k=1}^{2m} (X^k \circ \pi) \frac{\partial}{\partial \tilde{x}_k} - \sum_{i < j} \sum_{k < l} [y'_{kl} (G(\nabla_X S'_{kl}, S'_{ij}) \circ \pi) \frac{\partial}{\partial y'_{ij}} + y''_{kl} (G(\nabla_X S''_{kl}, S''_{ij}) \circ \pi) \frac{\partial}{\partial y''_{ij}}]$$

Let $a = (a', a'') \in \mathcal{A} \oplus \mathcal{A}$ and $p = \pi(a)$. Denote by $A(T_p M)$ the space of skew-symmetric endomorphisms of $(T_p M, g_p)$, the fibre of the bundle $\mathcal{A}$ at the point p . Then (9) implies that, under the standard identification of $T_a(A(T_p M) \oplus A(T_p M))$ with the vector space $A(T_p M) \oplus A(T_p M)$, we have

$$(10) \quad [X^h, Y^h]_a = [X, Y]_a^h + R(X, Y)a$$

where $R(X, Y)a = R(X, Y)a' + R(X, Y)a''$ is the curvature of the connection ∇ on the vector bundle $\mathcal{A} \oplus \mathcal{A}$ (for the curvature tensor we adopt the following definition: $R(X, Y) = \nabla_{[X, Y]} - [\nabla_X, \nabla_Y]$). Note also that (6) and (9) imply the well-known fact that

$$(11) \quad [V, X^h] \text{ is a vertical vector field.}$$

Notation. Let $J = (J_1, J_2) \in \mathcal{G}$ and $p = \pi(J)$. Take orthonormal bases $\{e'_1, \dots, e'_{2m}\}$, $\{e''_1, \dots, e''_{2m}\}$ of $T_p M$ such that $e'_{2l} = J_1 e'_{2l-1}$, $e''_{2l} = J_2 e''_{2l-1}$ for $l = 1, \dots, m$. Let $\{E'_i\}$, $\{E''_i\}$, $i = 1, \dots, 2m$, be orthonormal frames of TM near the point p such that

$$E'_i(p) = e'_i, \quad E''_i(p) = e''_i \quad \text{and} \quad \nabla E'_i|_p = 0, \quad \nabla E''_i|_p = 0, \quad i = 1, \dots, 2m.$$

For any (local) section $a = (a', a'')$ of $\mathcal{A} \oplus \mathcal{A}$, denote by $\tilde{a}$ the vertical vector field on $\mathcal{G}$ defined by

$$(12) \quad \tilde{a}_J = (a'_{\pi(J)} + J_1 \circ a'_{\pi(J)} \circ J_1, a''_{\pi(J)} + J_2 \circ a''_{\pi(J)} \circ J_2), \quad J = (J_1, J_2).$$

□

If S'_{ij} and S''_{ij} are the sections of bundle $\mathcal{A}$ defined by (4), then clearly $\nabla S'_{ij}|_p = \nabla S_{ij}|_p = 0$, $1 \leq i, j \leq 2m$. Note also that for every $J \in \mathcal{G}$ we can find sections $a_1, \dots, a_l$, $l = 2(m^2 - m)$, of $\mathcal{A} \oplus \mathcal{A}$ near the point $p = \pi(J)$ such that $\tilde{a}_1, \dots, \tilde{a}_l$ form a basis of the vertical vector space at each point in a neighbourhood of J .

The next lemma is a kind of folklore appearing in different contexts (cf, for example, [10, 4, 5]),

Lemma 1. *Let X be a vector field on M and let $a = (a', a'')$ be a section of $\mathcal{A} \oplus \mathcal{A}$ defined on a neighbourhood of the point $p = \pi(J)$. Then:*

$$[X^h, \tilde{a}]_J = (\widetilde{\nabla_X a})_J, \quad [X^h, \mathcal{J}^n \tilde{a}]_J = \mathcal{K}_n(\widetilde{\nabla_X a})_J.$$

Proof. Let $a'(E'_i) = \sum_{j=1}^{2m} a'_{ij} E'_j$, $a''(E''_i) = \sum_{j=1}^{2m} a''_{ij} E''_j$, $i = 1, \dots, 2m$. Then, in the local coordinates of $\mathcal{A} \oplus \mathcal{A}$ introduced above,

$$\tilde{a} = \sum_{i < j} [\tilde{a}'_{ij} \frac{\partial}{\partial y'_{ij}} + \tilde{a}''_{ij} \frac{\partial}{\partial y''_{ij}}]$$

where

$$(13) \quad \tilde{a}'_{ij} = a'_{ij} \circ \pi + \sum_{k,l=1}^{2m} y'_{ik} (a'_{kl} \circ \pi) y'_{lj}, \quad \tilde{a}''_{ij} = a''_{ij} \circ \pi + \sum_{k,l=1}^{2m} y''_{ik} (a''_{kl} \circ \pi) y''_{lj}.$$

By (9), we have for every vector field X on M near the point p

$$(14) \quad \begin{aligned} X^h_J &= \sum_{i=1}^{2m} X^i(p) \frac{\partial}{\partial x_i}(J), \\ [X^h, \frac{\partial}{\partial y'_{ij}}]_J &= [X^h, \frac{\partial}{\partial y''_{ij}}]_J = 0 \end{aligned}$$

since $\nabla S'_{ij}|_p = \nabla S''_{ij}|_p = 0$. Moreover,

$$(15) \quad (\nabla_{X_p} a')(E'_i) = \sum_{j=1}^{2m} X_p(a'_{ij}) E'_j, \quad (\nabla_{X_p} a'')(E''_i) = \sum_{j=1}^{2m} X_p(a''_{ij}) E''_j$$

since $\nabla E'_i|_p = \nabla E''_i|_p = 0$. Now, the lemma follows by a simple computation. □

Notation. Further on, we shall often make use of the standard isometric isomorphism $\mathcal{A} \cong \Lambda^2 TM$ that assigns to each $a \in A(T_p M)$ the 2-vector $a^\wedge$ for which

$$g(a^\wedge, X \wedge Y) = g(aX, Y), \quad X, Y \in T_p M,$$

where the metric on $\Lambda^2 TM$ is given by

$$g(X_1 \wedge X_2, X_3 \wedge X_4) = g(X_1, X_3)g(X_2, X_4) - g(X_1, X_4)g(X_2, X_3).$$

Under this isomorphism the connection on $\mathcal{A}$ induced by the Levi-Civita connection ∇ of (M, g) goes to the connection on $\Lambda^2 TM$ induced by ∇ . We shall use the same symbol ∇ for these connections.

Also, juxtaposition of endomorphisms will mean "composition". □

Lemma 2. ([4]) For every $a, b \in A(T_p M)$ and $X, Y \in T_p M$, we have

$$G(R(X, Y)a, b) = g(R([a, b]^\wedge)X, Y).$$

Corollary 1. If $J = (J_1, J_2) \in \mathcal{G}$, $X, Y \in \pi(J)$, and $V = (V_1, V_2) \in \mathcal{V}_J$

$$H_t(R(X, Y)J, V) = 2g(R(t_1(J_1 V_1)^\wedge + t_2(J_2 V_2)^\wedge)X, Y).$$

Proof. We have

$$\begin{aligned} H_t(R(X, Y)J, V) &= t_1 g(R([J_1, V_1]^\wedge)X, Y) + t_2 g(R([J_2, V_2]^\wedge)X, Y) \\ &= 2g(R(t_1(J_1 V_1)^\wedge + t_2(J_2 V_2)^\wedge)X, Y). \end{aligned}$$

□

Notation. The curvature operator of the connection ∇ will be denoted by $\mathcal{R}$. This is the self-adjoint endomorphism of $\Lambda^2 TM$ defined by

$$g(\mathcal{R}(X \wedge Y), Z \wedge T) = g(R(X, Y)Z, T), \quad X, Y, Z, T \in TM.$$

Also, we denote the Levi-Civita connection of the metric H_t on $\mathcal{G}$ by D . □

Lemma 3. ([3, 4]) If X, Y are vector fields on M and V is a vertical vector field on $Z(E)$, then:

(i) At any point $J \in \mathcal{G}$

$$(16) \quad (D_{X^h} Y^h)_J = (\nabla_X Y)_J^h + \frac{1}{2} R(X, Y)J.$$

(ii) The fibres of $\mathcal{G} \rightarrow M$ are totally geodesic submanifolds.

(iii) The vector $(D_V X^h)_J$ is horizontal and

$$(17) \quad H_t(D_V X^h, Y^h)_J = H_t(D_{X^h} V, Y^h)_J = -g(\mathcal{R}(t_1(J_1 V_1)^\wedge + t_2(J_2 V_2)^\wedge), X \wedge Y).$$

Proof. (i) Formula (16) follows from the Koszul formula and the fact that the Lie bracket of a vertical and a horizontal vector field is a vertical vector field

(ii) Set $S_{ij} = (S'_{ij}, S''_{ij})$ and, as in the preceding section, define

$$\begin{aligned} A_{r,s} &= \frac{1}{\sqrt{2}}(S_{2r-1,2s-1} - S_{2r,2s}), \quad B_{r,s} = \frac{1}{\sqrt{2}}(S_{2r-1,2s} + S_{2r,2s-1}), \\ r &= 1, \dots, m-1, \quad s = r+1, \dots, m. \end{aligned}$$

The vertical vector fields $\tilde{A}_{rs}, \tilde{B}_{rs}$ corresponding to the sections $A_{r,s}, B_{r,s}$ of $\mathcal{A} \oplus \mathcal{A}$ constitute a frame of the vertical bundle $\mathcal{V}$ of $\mathcal{G}$ in a neighbourhood of J . Let $V \in \mathcal{V}_J$. Take a section Q of $\mathcal{A} \oplus \mathcal{A}$ such that $Q(p) = V$ and $\nabla Q|_p = 0$ where $p = \pi(J)$. By Lemma 1, $[X^h, \tilde{A}_{rs}]_J = [X^h, \tilde{B}_{rs}]_J = [X^h, \tilde{Q}]_J = 0$ for every vector field X in a neighbourhood of p . It follows from the Koszul formula for the Levi-Civita connection that $(D_{\tilde{Q}} \tilde{A}_{rs})_J$ and $(D_{\tilde{Q}} \tilde{B}_{rs})_J$ are H_t -orthogonal to every horizontal vector X_J^h . Hence $D_V \tilde{A}_{rs}$ and $D_V \tilde{B}_{rs}$ are vertical tangent vectors of $\mathcal{G}$ at J . It follows that, for every vertical vector field W , $D_V W$ is a vertical vector field. Thus the fibres of $\mathcal{G}$ are totally geodesic submanifolds. This, of course, follows also from the Vilms theorem [16] (or [2]).

(iii) By (ii), $D_V X^h$ is orthogonal to every vertical vector field, thus $D_V X^h$ is a horizontal vector field. Hence $D_V X^h = \mathcal{H}D_{X^h} V$ since $[V, X^h]$ is vertical. Therefore

by (10) and Corollary 1

$$\begin{aligned} H_{\mathbf{t}}(D_V X^h, Y^h)_J &= H_{\mathbf{t}}(D_{X^h} V, Y^h)_J = -H_{\mathbf{t}}(V, D_{X^h} Y^h)_J = -\frac{1}{2} H_{\mathbf{t}}(R(X, Y)J, V) \\ &= -g(R(t_1(J_1 V_1)^\wedge + t_2(J_2 V_2)^\wedge)X, Y). \end{aligned}$$

□

Notation. Let $\Omega_{\mathbf{t},n}(A, B) = H_{\mathbf{t}}(\mathcal{J}^n A, B)$ be the fundamental 2-form of the almost Hermitian structure $(H_{\mathbf{t}}, \mathcal{J}^n)$ on $\mathcal{G}$, $n = 1, \dots, 4$. □

Proposition 1. *Let $J = (J_1, J_2) \in \mathcal{G}$, $X, Y, Z \in T_{\pi(J)}M$, and $U, V = (V_1, V_2), W \in \mathcal{V}_J$. Then*

$$(18) \quad (D_{Z^h} \Omega_{\mathbf{t},n})(X^h, Y^h)_J = 0$$

(19)

$$(D_V \Omega_{\mathbf{t},n})(X^h, Y^h)_J = g(V_1 X, Y) - g(\mathcal{R}(t_1(J_1 V_1)^\wedge + t_2(J_2 V_2)^\wedge), X \wedge J_1 Y + J_1 X \wedge Y).$$

where $(V_1, V_2) = V$.

(20)

$$\begin{aligned} (D_{Z^h} \Omega_{\mathbf{t},n})(X^h, V)_J &= (-1)^n g(\mathcal{R}(\pm t_1 V_1^\wedge + t_2 V_2^\wedge), Z \wedge X) \\ &\quad + g(\mathcal{R}(t_1(J_1 V_1)^\wedge + t_2(J_2 V_2)^\wedge), Z \wedge J_1 X), \end{aligned}$$

where the plus sign corresponds to $n = 1, 4$ and the minus sign to $n = 2, 3$.

$$(21) \quad (D_W \Omega_{\mathbf{t},n})(X^h, V) = 0, \quad (D_{Z^h} \Omega_{\mathbf{t},n})(U, V) = 0, \quad (D_W \Omega_{\mathbf{t},n})(U, V) = 0.$$

Proof. Extend X, Y to vector fields in a neighbourhood of $p = \pi(J)$ such that $\nabla X|_p = \nabla Y|_p = 0$. Note that $Z_J^h(y'_{kl}) = Z_J^h(y''_{kl}) = 0$, $k, l = 1, \dots, 2m$, by (9).

Then, by (8) and (16),

$$\begin{aligned} (D_{Z^h} \Omega_{\mathbf{t},n})(X^h, Y^h)_J &= Z_J^h H_{\mathbf{t}}(\mathcal{J}^n X^h, Y^h) \\ &\quad + H_{\mathbf{t}}(D_{Z^h} X^h, (J_1 Y)^h)_J - H_{\mathbf{t}}(D_{Z^h} Y^h, (J_1 X)^h)_J \\ &= \sum_{k=1}^{2m} \sum_{l=1}^{2m} [y'_{kl} Z((g(X, E'_k)g(Y, E'_l)) \circ \pi) = 0. \end{aligned}$$

Also, by (8) and (17),

$$\begin{aligned} (D_V \Omega_{\mathbf{t},n})(X^h, Y^h)_J &= \sum_{k=1}^{2m} \sum_{l=1}^{2m} V(y'_{kl}(g(X, E'_k)g(Y, E'_l)) \circ \pi) \\ &\quad + H_{\mathbf{t}}(D_V X^h, (J_1 Y)^h)_J - H_{\mathbf{t}}(D_V Y^h, (J_1 X)^h)_J \\ &= g(V_1 X, Y) - g(\mathcal{R}(t_1(J_1 V_1)^\wedge + t_2(J_2 V_2)^\wedge), X \wedge J_1 Y + J_1 X \wedge Y). \end{aligned}$$

Next

$$\begin{aligned} (D_{Z^h} \Omega_{\mathbf{t},n})(X^h, V)_J &= H_{\mathbf{t}}(D_{Z^h} X^h, \mathcal{J}^n V) - H_{\mathbf{t}}((J_1 X)^h, D_{Z^h} V) \\ &= \frac{1}{2} H_{\mathbf{t}}(R(Z, X)J, \mathcal{K}_n V) + g(\mathcal{R}(t_1(J_1 V_1)^\wedge + t_2(J_2 V_2)^\wedge), Z \wedge J_1 X) \end{aligned}$$

and (20) follows from Corollary 1

Next, by Lemma 3

$$(D_W \Omega_{\mathbf{t},n})(X^h, V) = 0.$$

In order to prove the second identity in (21), take sections $a = (a', a'')$ and $b = (b', b'')$ of $\mathcal{A} \oplus \mathcal{A}$ such that $a(p) = U$, $b(p) = V$ and $\nabla a|_p = \nabla b|_p = 0$. Denote by

$\tilde{a}$ and $\tilde{b}$ the vertical vector fields on $\mathcal{G}$ determined by a and b via (12). Then, by Lemma 3 (iii) and Lemma 1

$$\begin{aligned} (D_{Z^h} \Omega_{\mathbf{t},n})(U, V) &= \frac{1}{4} [Z_J^h H_{\mathbf{t}}(\mathcal{J}^n \tilde{a}, \tilde{b}) + H_{\mathbf{t}}([Z^h, \tilde{a}], \mathcal{J}^n \tilde{b})_J - H_{\mathbf{t}}(\mathcal{J}^n \tilde{a}, [Z^h, \tilde{b}])_J] \\ &= \frac{1}{4} Z_J^h H_{\mathbf{t}}(\mathcal{J}^n \tilde{a}, \tilde{b}) \end{aligned}$$

Set $a'(E'_i) = \sum_{j=1}^{2m} a'_{ij} E'_j$, $a''(E''_i) = \sum_{j=1}^{2m} a''_{ij} E''_j$, $i = 1, \dots, 2m$, and define b'_{ij} , b''_{ij} in a similar way. We have

$$\tilde{a} = \sum_{i < j} [\tilde{a}'_{ij} \frac{\partial}{\partial y'_{ij}} + \tilde{a}''_{ij} \frac{\partial}{\partial y''_{ij}}], \quad \tilde{b} = \sum_{i < j} [\tilde{b}'_{ij} \frac{\partial}{\partial y'_{ij}} + \tilde{b}''_{ij} \frac{\partial}{\partial y''_{ij}}],$$

where $\tilde{a}'_{ij}, \tilde{a}''_{ij}, \tilde{b}'_{ij}, \tilde{b}''_{ij}$ are given by (13). Then $Z_J^h(\tilde{a}'_{ij}) = Z_J^h(\tilde{a}''_{ij}) = 0$, $Z_J^h(\tilde{b}'_{ij}) = Z_J^h(\tilde{b}''_{ij}) = 0$. Hence, by (7),

$$Z_J^h H_{\mathbf{t}}(\mathcal{J}^n \tilde{a}, \tilde{b}) = (-1)^{n+1} \sum_{j < k} \sum_{s=1}^{2m} Z_J^h [\pm \tilde{a}'_{js} y'_{sk} \tilde{b}'_{jk} + \tilde{a}''_{js} y''_{sk} \tilde{b}''_{jk}] = 0.$$

This proves the second identity in (21). The third one follows from the fact that the fibres of $\mathcal{G}$ are totally geodesic and Kähler. $\square$

Corollary 2. Let $J = (J_1, J_2) \in \mathcal{G}$, $X, Y, Z \in T_{\pi(J)} M$, and $U, V = (V_1, V_2), W \in \mathcal{V}_J$. Then

$$d\Omega_{\mathbf{t},n}(X^h, Y^h, Z^h) = 0.$$

$$d\Omega_{\mathbf{t},n}(X^h, Y^h, V)_J = g(V_1 X, Y) + 2(-1)^n g(\mathcal{R}(\pm t_1 V_1^\wedge + t_2 V_2^\wedge), X \wedge Y),$$

where the plus sign corresponds to $n = 1, 4$ and the minus sign to $n = 2, 3$.

$$d\Omega_{\mathbf{t},n}(X^h, V, W) = 0, \quad d\Omega_{\mathbf{t},n}(U, V, W) = 0$$

Corollary 3. Let $J = (J_1, J_2) \in \mathcal{G}$, $X \in T_{\pi(J)} M$, and $V = (V_1, V_2) \in \mathcal{V}_J$. Then

$$(\delta\Omega_{\mathbf{t},n})(X^h) = 0, \quad (\delta\Omega_{\mathbf{t},n})(V) = -2g(\mathcal{R}(t_1(J_1 V_1)^\wedge + t_2(J_2 V_2)^\wedge), J_1^\wedge),$$

Denote by $\mathcal{N}_n$ the Nijenhuis tensor of the almost complex structure $\mathcal{J}^n$:

$$\mathcal{N}_n(A, B) = -[A, B] + [\mathcal{J}^n A, \mathcal{J}^n B] - \mathcal{J}^n[\mathcal{J}^n A, B] - \mathcal{J}^n[A, \mathcal{J}^n B], \quad A, B \in T\mathcal{G}.$$

The Nejenhuis tensor is related to the covariant derivatives of the 2-form $\Omega_{\mathbf{t},n}$ by the following well-known identity

$$\begin{aligned} H_{\mathbf{t}}(\mathcal{N}_n(A, B), C) &= (D_A \Omega_{\mathbf{t},n})(\mathcal{J}^n B, C) - (D_B \Omega_{\mathbf{t},n})(\mathcal{J}^n A, C) \\ &\quad + (D_{\mathcal{J}^n A})(B, C) - (D_{\mathcal{J}^n B})(A, C) \end{aligned}$$

This and Proposition 1 imply:

Corollary 4. Let $J = (J_1, J_2) \in \mathcal{G}$, $X, Y, Z \in T_{\pi(J)} M$, and $V = (V_1, V_2), W \in \mathcal{V}_J$. Then

$$H_{\mathbf{t},n}(\mathcal{N}_n(X^h, Y^h), Z^h) = 0.$$

$$\begin{aligned} H_{\mathbf{t},n}(\mathcal{N}_n(X^h, Y^h), V)_J &= 2(-1)^n (g(\mathcal{R}(\pm t_1 V_1^\wedge + t_2 V_2^\wedge), X \wedge J_1 Y + J_1 X \wedge Y) \\ &\quad - 2g(\mathcal{R}(t_1(J_1 V_1)^\wedge + t_2(J_2 V_2)^\wedge), X \wedge Y - J_1 X \wedge J_1 Y)) \end{aligned}$$

where the plus sign corresponds to $n = 1, 4$ and the minus sign to $n = 2, 3$.

$$\begin{aligned} H_{\mathbf{t},n}(\mathcal{N}_n(X^h, V), Y^h) &= 0 \quad \text{for } n = 1, 2, \\ H_{\mathbf{t},n}(\mathcal{N}_n(X^h, V), Y^h)_J &= 2g(J_1 V_1 X, Y) \quad \text{for } n = 3, 4. \\ H_{\mathbf{t},n}(\mathcal{N}_n(X^h, V), W) &= 0, \quad \mathcal{N}_n(V.W) = 0. \end{aligned}$$

4. THE GRAY-HERVELLA CLASSES

Let $\tilde{\Omega}_{t,k}$, $t > 0$, $k = 1, 2$, be the fundamental 2-form of the almost Hermitian structure $(h_t, \mathcal{J}_k)$ on the twistor space $\mathcal{Z}$. Let $\tilde{D}$ be the Levi-Civita connection of h_t . Considering $\mathcal{Z}$ as a subbundle of $\mathcal{A}$, computations similar to those in the proof of Proposition 1 give formulas for the covariant derivatives $\tilde{D}\tilde{\Omega}_{t,k}$ which show the following relation with the covariant derivative $D\Omega_{\mathbf{t},n}$.

Let $J = (J_1, J_2) \in \mathcal{G}$. Recall that, considering $\mathcal{G}$ as a subbundle of $\mathcal{A} \oplus \mathcal{A}$, the horizontal space of $\mathcal{G}$ at J is the horizontal space of the Whitney sum $\mathcal{A} \oplus \mathcal{A}$ at J , the latter being the direct sum of the horizontal spaces of the vector bundle $\mathcal{A}$ at J_1 and J_2 . Note also that the horizontal space of $\mathcal{A}$ at J_k , $k = 1, 2$, is the horizontal space of $\mathcal{Z}$ at J_k . Then if $X \in T_{\pi(J)}M$, denoting the horizontal lift of X to $T_{J_k}\mathcal{Z}$ by $X_{J_k}^{\tilde{h}}$, we have $X_J^h = (X_{J_1}^{\tilde{h}}, X_{J_2}^{\tilde{h}})$. Let $X, Y, Z \in T_{\pi(J)}M$, $U_1, V_1, W_1 \in \mathcal{V}_{J_1}$ (the vertical subspace of $T_{J_1}\mathcal{Z}$). Then for $U = (U_1, 0), V = (V_1, 0), W = (W_1, 0) \in \mathcal{V}_J$, for every $\mathbf{t} = (t_1, t_2)$, $t_1, t_2 > 0$, and for $n = 1, 2$,

$$\begin{aligned} (\tilde{D}_{Z_{J_1}^{\tilde{h}} + W_1} \tilde{\Omega}_{t_1,1})(X_{J_1}^{\tilde{h}} + U_1, Y_{J_1}^{\tilde{h}} + V_1) &= (D_{Z_J^h + W} \Omega_{\mathbf{t},n})(X_J^h + U, Y_J^h + V), \\ d\tilde{\Omega}_{t_1,1}(X_{J_1}^{\tilde{h}} + U_1, Y_{J_1}^{\tilde{h}} + V_1, Z_{J_1}^{\tilde{h}} + W_1) &= d\Omega_{\mathbf{t},n}(X_J^h + U, Y_J^h + V, Z_J^h + W), \\ \delta\tilde{\Omega}_{t_1,1}(X_{J_1}^{\tilde{h}} + U_1) &= \delta\Omega_{\mathbf{t},n}(X_J^h + U), \end{aligned}$$

Similar identities holds for the covariant derivatives, differentials and codifferentials of $\tilde{\Omega}_{t_1,2}$ and $\Omega_{\mathbf{t},n}$ for $n = 3, 4$. Note also that

$$h_{t_1}(X_{J_1}^{\tilde{h}} + U_1, Y_{J_1}^{\tilde{h}} + V_1) = H_{(t_1, t_2)}(X_J^h + U, Y_J^h + V).$$

These relations and the defining identities for the Gray-Hervella classes [11, p. 41] imply the following.

Proposition 2. *If one of the almost Hermitian structures $(H_{\mathbf{t}}, \mathcal{J}^1)$ and $(H_{\mathbf{t}}, \mathcal{J}^2)$ on $\mathcal{G}$ belongs to a (non-trivial) Gray-Hervella class, then the almost Hermitian structure $(h_{t_1}, \mathcal{J}_1)$ on $\mathcal{Z}$ belongs to the same class.*

If one of $(H_{\mathbf{t}}, \mathcal{J}^3)$ and $(H_{\mathbf{t}}, \mathcal{J}^4)$ is in a Gray-Hervella class, $(h_{t_1}, \mathcal{J}_2)$ is in the same class.

Suppose that the manifold M is oriented and denote by $\mathcal{Z}_{\pm}$ the bundle over M whose sections are the almost complex structures compatible with the metric and $\pm$ the orientation of M . The bundles $\mathcal{Z}_+$ and $\mathcal{Z}_-$ are usually called the positive and the negative twistor space of (M, g) , respectively. The spaces $\mathcal{Z}_{\pm}$ are disjoint open subsets of the twistor space $\mathcal{Z}$. If M is connected and oriented, they are the connected components of the manifold $\mathcal{Z}$. Correspondingly, $\mathcal{G} = \mathcal{Z} \times_M \mathcal{Z}$ has four connected components, the product bundles $\mathcal{Z}_+ \times_M \mathcal{Z}_+$, $\mathcal{Z}_+ \times_M \mathcal{Z}_-$, $\mathcal{Z}_- \times_M \mathcal{Z}_+$, $\mathcal{Z}_- \times_M \mathcal{Z}_-$.

Clearly, Proposition 2 holds for the restrictions of $(H_{\mathbf{t}}, \mathcal{J}^n)$ to $\mathcal{Z}_+ \times_M \mathcal{Z}_+$ and $\mathcal{Z}_+ \times_M \mathcal{Z}_-$ (resp. $\mathcal{Z}_- \times_M \mathcal{Z}_+$ and $\mathcal{Z}_- \times_M \mathcal{Z}_-$) and the restriction of $(h_{t_1}, \mathcal{J}_k)$ to $\mathcal{Z}_+$ (resp. $\mathcal{Z}_-$).

Recall that the curvature operator has the decomposition [2, G (1.116)].

$$\mathcal{R} = \frac{s}{n(n-1)} Id + \mathcal{B} + \mathcal{W}$$

where $s = s_M$ is the scalar curvature, the operator $\mathcal{B}$ corresponds to the traceless Ricci tensor, and $\mathcal{W}$ corresponds to the Weyl conformal tensor. As is well-known, if M is oriented and four-dimensional, the operator $\mathcal{W}$ has an extra decomposition which can be described by means of the Hodge star operator. In dimension four, the Hodge operator acts on $\Lambda^2 TM$ as an involution and we have the orthogonal decomposition

$$\Lambda^2 TM = \Lambda_+^2 TM \oplus \Lambda_-^2 TM$$

where $\Lambda_{\pm}^2 TM$ are the subbundles corresponding to the ± 1 -eigenvalues of the Hodge operator. Denoting the restriction of $\mathcal{W}$ to $\Lambda_{\pm}^2 TM$ by $\mathcal{W}_{\pm}$, we obtain the decomposition ([15], [2, G sec. 1.128])

$$(22) \quad \mathcal{R} = \frac{s}{12} Id + \mathcal{B} + \mathcal{W}_+ + \mathcal{W}_-.$$

Note that $\mathcal{B}$ sends $\Lambda_{\pm}^2 TM$ to $\Lambda_{\mp}^2 TM$, $\mathcal{W}_{\pm}$ sends $\Lambda_{\pm}^2 TM$ to $\Lambda_{\pm}^2 TM$, and $\mathcal{W}_{\pm}|_{\Lambda_{\mp}^2 TM} = 0$. Note also that changing the orientation of M interchanges the roles of $\Lambda_+^2 TM$ and $\Lambda_-^2 TM$, respectively the roles of $\mathcal{W}_+$ and $\mathcal{W}_-$.

The manifold (M, g) is Einstein exactly when $\mathcal{B} = 0$. It is called anti-self-dual (self-dual) when $\mathcal{W}_+ = 0$ ($\mathcal{W}_- = 0$).

Recall that by the Atiyah-Hitchin-Singer theorem [1] the almost complex structure $\mathcal{J}_1$ restricted to $\mathcal{Z}_+$ (resp. $\mathcal{Z}_-$) is integrable if and only if the base manifold (M, g) is anti-self-dual (resp. self-dual). In contrast, the almost complex structure $\mathcal{J}_2$ is never integrable by a result of Eells-Salamon [8].

The isometry $\mathcal{A} \cong \Lambda^2 TM$, $a \rightarrow a^\wedge$, identifies the bundle $\mathcal{Z}_{\pm}$ with the sphere subbundle of $\Lambda_{\pm}^2 TM$ of radius $\sqrt{2}$. Under this isomorphism, if $J \in \mathcal{Z}_{\pm}$, the vertical space $\mathcal{V}_J$ at J is identified with the orthogonal complement $(\mathbb{R}J^\wedge)^\perp$ of $\mathbb{R}J^\wedge$ in $\Lambda_{\pm}^2 TM$.

Convention. Henceforth, we assume that M is oriented and of dimension four. We shall freely identify skew-symmetric endomorphisms with 2-forms and shall use the same notation for both of them. $\square$

Notation. It is convenient to set $\mathcal{G}_{++} = \mathcal{Z}_+ \times_M \mathcal{Z}_+$, $\mathcal{G}_{+-} = \mathcal{Z}_+ \times_M \mathcal{Z}_-$, and similar for $\mathcal{G}_{--}$, $\mathcal{G}_{-+}$. $\square$

Changing the orientation of M interchanges the roles of $\mathcal{G}_{++}$ and $\mathcal{G}_{--}$, as well as $\mathcal{G}_{+-}$ and $\mathcal{G}_{-+}$. Hence it is enough to consider only $\mathcal{G}_{++}$ and $\mathcal{G}_{+-}$.

According to [14, Theorem 4.6] the possible Gray-Hervella classes for $(h_{t_1}, \mathcal{J}_1)$ considered on $\mathcal{Z}_+$ are $\mathcal{K}$ and $\mathcal{W}_3$. Thus Proposition 2 implies the following

Corollary 5. *The only possible Gray-Hervella classes for $(H_{\mathbf{t}}, \mathcal{J}^1)$ and $(H_{\mathbf{t}}, \mathcal{J}^2)$ restricted to $\mathcal{G}_{++} = \mathcal{Z}_+ \times_M \mathcal{Z}_+$ or $\mathcal{G}_{+-} = \mathcal{Z}_+ \times_M \mathcal{Z}_-$ are $\mathcal{K}$ and $\mathcal{W}_3$.*

Remark. The class $\mathcal{K}$ consists of Kähler structures; $\mathcal{W}_3$ is the class of Hermitian semi-Kähler structures (for the defining condition of this class see the proof of the Theorem 2 below).

Theorem 1.

- (a) $(H_t, \mathcal{J}^n)$, $n = 1$ or 2 , restricted to $\mathcal{G}_{++}$ does not belong to the class $\mathcal{K}$.
- (b) $(H_t, \mathcal{J}^n)$, $n = 1$ or 2 , restricted to $\mathcal{G}_{+-}$ belongs to the class $\mathcal{K}$ if and only if (M, g) is Einstein with positive scalar curvature s , anti-self-dual, $\mathcal{W}_-(\tau) = -\frac{s}{12}\tau$ for $\tau \in \Lambda_-^2 TM$, and $t_1 = \frac{6}{s}$.

Proof. Suppose that one of the structures $(H_t, \mathcal{J}^n)$, $n = 1$ or 2 , restricted to $\mathcal{Z}_+ \times_M \mathcal{Z}_+$ or $\mathcal{Z}_+ \times_M \mathcal{Z}_-$ belongs to the class $\mathcal{K}$. Then so does the restriction of $(h_{t_1}, \mathcal{J}_1)$ to $\mathcal{Z}_+$. It is well-known that $(h_{t_1}, \mathcal{J}_1)$ is Kähler on $\mathcal{Z}_+$ if and only if the base manifold (M, g) is Einstein with positive scalar curvature, anti-self-dual, and $t_1 = \frac{6}{s}$ [9] (see also [14], but be aware that the metric on $\Lambda^2 TM$ used there is one half of the metric here and in [9], hence the curvature operator $\mathcal{R}$ is one half of the curvature operator in [14]). Then by (19)

$$(23) \quad g(\mathcal{R}(J_2 V_2), X \wedge J_1 Y + J_1 X \wedge Y) = 0$$

for every $J = (J_1, J_2) \in \mathcal{Z}_+ \times_M \mathcal{Z}_+$ or $\mathcal{Z}_+ \times_M \mathcal{Z}_-$, $X, Y \in T_{\pi(J)} M$, $V_2 \in \mathcal{V}_{J_2}$, where $\mathcal{R}(J_2 V_2) = \frac{s}{12} J_2 V_2 + \mathcal{W}_-(J_2 V_2)$. Note that the 2-vector $X \wedge J_1 Y + J_1 X \wedge Y \in \Lambda_+^2 T_{\pi(J)} M$. This follows from the easily verifiable fact that $X \wedge J_1 Y + J_1 X \wedge Y$ is orthogonal to $\Lambda_-^2 T_p M$, $p = \pi(J)$, since J_1 commutes with every endomorphism of $T_p M$ lying in $\Lambda_-^2 T_p M$.

- (a). If $J_2 \in \Lambda_+^2 TM$, identity (23) takes the form

$$g(J_2 V_2, X \wedge J_1 Y + J_1 X \wedge Y) = 0.$$

The latter identity is not satisfied, for example, for $J = \sqrt{2}(s_1^+, s_2^+)$, $V_2 = s_1^+$, $(X, Y) = (E_1, E_3)$ where $E_1, \dots, E_4$ is an oriented orthonormal basis of $T_p M$ and $\{s_i^+\}$ are defined via (3). Therefore none of the restriction of $(H_t, \mathcal{J}^1)$ and $(H_t, \mathcal{J}^2)$ to $\mathcal{G}_{++}$ is a Kähler structure.

- (b). If $J_2 \in \Lambda_-^2 TM$, identity (23) is automatically satisfied since $\mathcal{R}(J V_2) \in \Lambda_-^2 T_p M$. We also have by (20)

$$g(\mathcal{R}(V_2), Z \wedge X) = g(\mathcal{R}(J_2 V_2), Z \wedge J_1 X).$$

This implies

$$g(\mathcal{R}(V_2), Z \wedge X + J_1 Z \wedge J_1 X) = g(\mathcal{R}(J_2 V_2), Z \wedge J_1 X - J_1 Z \wedge X),$$

$$g(\mathcal{R}(V_2), Z \wedge X - J_2 Z \wedge J_2 X) = g(\mathcal{R}(J_2 V_2), Z \wedge J_1 X - J_2 Z \wedge J_1 J_2 X).$$

Set $J_1 = \sqrt{2}s_1^+$, $(Z, X) = (E_1, E_3)$ in the first of the latter identities and also $(J_1, J_2) = \sqrt{2}(-s_1^+, s_1^-)$, $(Z, X) = (E_1, E_3)$. Thus, we obtain

$$g(\mathcal{R}(V_2), s_2^-) = g(\mathcal{R}(J_2 V_2), s_3^-), \quad g(\mathcal{R}(V_2), s_2^-) = -g(\mathcal{R}(J_2 V_2), s_3^-),$$

where $J_2 = \sqrt{2}s_1^-$ and $V_2 \perp s_1^-$. Hence

$$g(\mathcal{R}(s_2^-), s_2^-) = g(\mathcal{R}(s_3^-), s_2^-) = 0.$$

Replacing the basis E_1, E_2, E_3, E_4 with the bases E_1, E_3, E_4, E_2 and E_1, E_4, E_2, E_3 , we see that

$$g(\mathcal{R}(s_3^-), s_3^-) = g(\mathcal{R}(s_1^-), s_3^-) = 0, \quad g(\mathcal{R}(s_1^-), s_1^-) = g(\mathcal{R}(s_2^-), s_1^-) = 0.$$

It follows that $\mathcal{R}|\Lambda_-^2 TM = 0$. Thus, if one of the restrictions of the structures $(H_t, \mathcal{J}^1)$ and $(H_t, \mathcal{J}^2)$ to $\mathcal{Z}_+ \times_M \mathcal{Z}_-$ is Kähler, then (M, g) is Einstein with positive scalar curvature s , anti-self-dual, $\mathcal{W}_- = -\frac{s}{12}Id|\Lambda_-^2 TM$, and $t_1 = \frac{6}{s}$.

Conversely, if (M, g) satisfies these conditions, it follows easily from Proposition 1 that the restrictions of $(H_t, \mathcal{J}^1)$ and $(H_t, \mathcal{J}^2)$ to $\mathcal{G}_{+-}$ are Kähler structures. $\square$

Theorem 2.

- (a) $(H_t, \mathcal{J}^n)$, $n = 1$ or 2 , restricted to $\mathcal{G}_{++}$ is in the class $\mathcal{W}_3$ if and only if (M, g) is anti-self-dual and scalar flat.
- (b) $(H_t, \mathcal{J}^n)$, $n = 1$ or 2 , restricted to $\mathcal{G}_{+-}$ is in the class $\mathcal{W}_3$ if and only if (M, g) is anti-self-dual.

Proof. (a). Suppose that one of the restrictions of $(H_t, \mathcal{J}^1)$ and $(H_t, \mathcal{J}^2)$ to $\mathcal{Z}_+ \times_M \mathcal{Z}_+$ belongs to the class $\mathcal{W}_3$. Recall that the defining conditions for this class are: (i) the Nijenhuis tensor of the almost complex structure vanishes; (ii) the codifferential of the fundamental 2-form vanishes. By Proposition 2, the almost Hermitian structure $(h_{t_1}, \mathcal{J}_1)$ on $\mathcal{Z}_+$ is of class $\mathcal{W}_3$. In particular, the Atiyah-Hitchin-Singer structure $\mathcal{J}_1$ on $\mathcal{Z}_+$ is integrable. Therefore (M, g) is anti-self-dual. Then, by Corollary 3, condition (ii) takes the form $g(\mathcal{R}(J_2 V_2), J_1) = 0$. Thus

$$\frac{s}{12}g(J_2 V_2, J_1) = 0.$$

for every $J_1, J_2 \in \mathcal{Z}_+$ and every $V_2 \in \mathcal{V}_{J_2}$. This implies $s = 0$.

Conversely, if (M, g) is anti-self-dual and scalar flat, then conditions (i) and (ii) are fulfilled by Corollaries 4 and 3.

(b). Suppose that one of the restrictions of $(H_t, \mathcal{J}^1)$ and $(H_t, \mathcal{J}^2)$ to $\mathcal{Z}_+ \times_M \mathcal{Z}_-$ belongs to the class $\mathcal{W}_3$. Then (M, g) is anti-self-dual. Conversely, if (M, g) is anti-self dual, $\mathcal{R}|\Lambda_-^2 TM$ takes its values in $\Lambda_-^2 TM$. Then conditions (i) and (ii) are satisfied by Corollaries 4 and 3 since $X \wedge J_1 Y + J_1 X \wedge Y$ and $X \wedge Y - J_1 X \wedge J_1 Y$ lies in $\Lambda_+^2 TM$. $\square$

Proposition 2 and [14, Theorem 4.7] imply the following

Corollary 6. *The only possible Gray-Hervella classes for $(H_t, \mathcal{J}^3)$ and $(H_t, \mathcal{J}^4)$ restricted to $\mathcal{G}_{++} = \mathcal{Z}_+ \times \mathcal{Z}_+$ or $\mathcal{G}_{+-} = \mathcal{Z}_+ \times \mathcal{Z}_-$ are $\mathcal{W}_1 \oplus \mathcal{W}_2 \oplus \mathcal{W}_3$, $\mathcal{W}_1 \oplus \mathcal{W}_2$, $\mathcal{W}_1 \oplus \mathcal{W}_3$, $\mathcal{W}_2 \oplus \mathcal{W}_3$, $\mathcal{W}_1$, $\mathcal{W}_2$.*

Remark. $\mathcal{W}_1 \oplus \mathcal{W}_2 \oplus \mathcal{W}_3$ is the class of semi-Kähler almost Hermitian structures, $\mathcal{W}_1 \oplus \mathcal{W}_2$ of the quasi-Kähler ones, $\mathcal{W}_1$ is the class of nearly-Kähler structures, $\mathcal{W}_2$ of almost-Kähler ones. The defining conditions of these classes are given in the proofs of the next claims. .

Theorem 3.

- (a) $(H_t, \mathcal{J}^n)$, $n = 3$ or 4 , restricted to $\mathcal{G}_{++}$ is in the class $\mathcal{W}_1 \oplus \mathcal{W}_2 \oplus \mathcal{W}_3$, if and only if (M, g) is anti-self-dual and scalar flat.
- (b) $(H_t, \mathcal{J}^n)$, $n = 3$ or 4 , restricted to $\mathcal{G}_{+-}$ is in the class $\mathcal{W}_1 \oplus \mathcal{W}_2 \oplus \mathcal{W}_3$ if and only if (M, g) is anti-self-dual.

Proof. The class $\mathcal{W}_1 \oplus \mathcal{W}_2 \oplus \mathcal{W}_3$ is defined by the requirement that the codifferential of the fundamental 2-form vanishes. For $n = 3$ or 4 , if one of the restrictions of $(H_t, \mathcal{J}^n)$ to $\mathcal{Z}_+ \times_M \mathcal{Z}_+$ belongs to this class, then (M, g) is anti-self-dual by

Proposition 2 and [14, Theorem 4.7 (i)]. Then, as we have seen in the proof of Theorem 2: (a) (M, g) is scalar flat when we consider the defining condition on $\mathcal{Z}_+ \times_M \mathcal{Z}_+$, (b) no other restrictions in the case of $\mathcal{Z}_+ \times_M \mathcal{Z}_-$. This proves the "if" part of the theorem. The converse statement follows immediately from Corollary 3. $\square$

Theorem 4.

- (a) $(H_t, \mathcal{J}^n)$, $n = 3$ or 4 , restricted to $\mathcal{G}_{++}$ is in the class $\mathcal{W}_1 \oplus \mathcal{W}_2$ if and only if (M, g) is anti-self-dual and Ricci flat
- (b) $(H_t, \mathcal{J}^n)$, $n = 3$ or 4 , restricted to $\mathcal{G}_{+-}$ is in the class $\mathcal{W}_1 \oplus \mathcal{W}_2$ if and only if (M, g) is Einstein, anti-self-dual, and $\mathcal{W}_-(\tau) = -\frac{s}{12}\tau$ for $\tau \in \Lambda_-^2 TM$.

Proof. Suppose that one of the restrictions of $(H_t, \mathcal{J}^3)$ and $(H_t, \mathcal{J}^4)$ to $\mathcal{Z}_+ \times_M \mathcal{Z}_+$ or $\mathcal{Z}_+ \times_M \mathcal{Z}_-$ belongs to $\mathcal{W}_1 \oplus \mathcal{W}_2$. The defining condition of this class reads as

$$(D_A \Omega_{t,n})(B, C) + (D_{\mathcal{J}^n A} \Omega_{t,n})(\mathcal{J}^n B, C) = 0, \quad A, B, C \in T\mathcal{G}, \quad n = 3, 4.$$

By Proposition 2 and [14, Theorem 4.7 (ii)], (M, g) is Einstein and anti-self-dual. Under these conditions, it is easy to see by means of Proposition 1 that the defining condition is equivalent to

$$(24) \quad \begin{aligned} g(\mathcal{R}(V_2), X \wedge Y) + g(\mathcal{R}(J_2 V_2), X \wedge J_1 Y) &= 0, \quad \text{for } n = 3, \\ g(\mathcal{R}(V_2), X \wedge Y) - g(\mathcal{R}(J_2 V_2), J_1 X \wedge Y) &= 0, \quad \text{for } n = 4. \end{aligned}$$

- (a). If $V_2 \in \Lambda_+^2 TM$, the latter identities become

$$\frac{s}{12}[g(V_2 X, Y) + g((J_2 V_2)X, J_1 Y)] = 0, \quad \frac{s}{12}[g(V_2 X, Y) + g((J_2 V_2)Y, J_1 X)] = 0.$$

Setting $J_1 = \sqrt{2}s_1^+$, $J_2 = \sqrt{2}s_2^+$, $V_2 = s_1^+$, $(X, Y) = (E_3, E_3)$ for an oriented orthonormal basis $E_1, \dots, E_4$, we see that $s = 0$.

- (b). If $V_2 \in \Lambda_-^2 TM$, identities (24) with $J_1 = \sqrt{2}s_1^+$ and $(X, Y) = (E_1, E_2)$ imply $g(\mathcal{R}(V_2), E_1 \wedge E_2) = 0$. It follows that $\mathcal{R}(V_2) = 0$ for every $V_2 \in \Lambda_-^2 TM$.

This proves the necessity of the curvature conditions stated in the theorem. Their sufficiency is an easy consequence of Proposition 1. $\square$

Theorem 5.

- (a) $(H_t, \mathcal{J}^n)$, $n = 3$ or 4 , restricted to $\mathcal{G}_{++}$ does not belong to the class $\mathcal{W}_1 \oplus \mathcal{W}_3$.
- (b) $(H_t, \mathcal{J}^n)$, $n = 3$ or 4 , restricted to $\mathcal{G}_{+-}$ belongs to the class $\mathcal{W}_1 \oplus \mathcal{W}_3$ if and only if (M, g) is Einstein with positive scalar curvature s , anti-self-dual, and $t_1 = \frac{3}{s}$.

Proof. The defining conditions of the class $\mathcal{W}_1 \oplus \mathcal{W}_3$ in the case of $(H_t, \mathcal{J}^n)$ read as (25)

$$(D_A \Omega_{t,n})(A, C) + (D_{\mathcal{J}^n A} \Omega_{t,n})(\mathcal{J}^n A, C) = 0, \quad \delta \Omega_{t,n} = 0, \quad A, C \in T\mathcal{G}, \quad n = 3, 4.$$

If one of the restrictions of $(H_t, \mathcal{J}^n)$ to $\mathcal{Z}_+ \times_M \mathcal{Z}_+$ or $\mathcal{Z}_+ \times_M \mathcal{Z}_-$ belongs to $\mathcal{W}_1 \oplus \mathcal{W}_3$, then (M, g) is anti-self dual with positive scalar curvature $s = \frac{3}{t_1}$ by Proposition 2

and [14, Theorem 4.7 (iii)]. Then identities (25) are equivalent to

$$(26) \quad (-1)^n g(\mathcal{R}(V_2), X \wedge Y - J_1 X \wedge J_1 Y) + g(\mathcal{R}(J_2 V_2), X \wedge J_1 Y + J_1 X \wedge Y) = 0, \\ g(\mathcal{R}(J_2 V_2), J_1) = 0$$

for every $J = (J_1, J_2) \in \mathcal{G}_{++}$ or $\mathcal{G}_{+-}$, $X, Y \in T_{\pi(J)}M$, $V_2 \in \mathcal{V}_{J_2}$.

For $J \in \mathcal{G}_{++}$ and $V_2 \in \Lambda_+^2 TM$, the second equation of (26) becomes

$$\frac{s}{12} g(J_2 V_2, J_1) = 0.$$

This implies $s = 0$, a contradiction. This proves (a).

If $J \in \mathcal{G}_{+-}$ and $V_2 \in \Lambda_-^2 TM$, then the second equation of (26) becomes

$$\frac{s}{12} g(\mathcal{B}(J_2 V_2), J_1) = 0.$$

This is equivalent to $\mathcal{B} = 0$. The 2-vectors $X \wedge Y - J_1 X \wedge J_1 Y$ and $X \wedge J_1 Y + J_1 X \wedge Y$ in the first identity of (26) lie in $\Lambda_+^2 TM$. Hence this identity is always satisfied when $\mathcal{B} = 0$ since in this case $\mathcal{R}$ preserves $\Lambda_-^2 TM$. This proves (b). $\square$

Theorem 6.

- (a) $(H_t, \mathcal{J}^n)$, $n = 3$ or 4 , restricted to $\mathcal{G}_{++}$ does not belong to the class $\mathcal{W}_2 \oplus \mathcal{W}_3$.
- (b) $(H_t, \mathcal{J}^n)$, $n = 3$ or 4 , restricted to $\mathcal{G}_{+-}$ belongs to the class $\mathcal{W}_2 \oplus \mathcal{W}_3$ if and only if (M, g) is Einstein with negative scalar curvature s , anti-self-dual, and $t_1 = -\frac{6}{s}$.

Proof. The belonging conditions for the class $\mathcal{W}_2 \oplus \mathcal{W}_3$ are

$$(27) \quad \mathfrak{S}_{A,B,C} \{ (D_A \Omega_{t,n})(A, C) - (D_{\mathcal{J}^n A} \Omega_{t,n})(\mathcal{J}^n A, C) \} = 0, \quad \delta \Omega_{t,n} = 0, \quad A, C \in T\mathcal{G}, \quad n = 3, 4.$$

If one of the restrictions of $(H_t, \mathcal{J}^n)$ to $\mathcal{Z}_+ \times_M \mathcal{Z}_+$ or $\mathcal{Z}_+ \times_M \mathcal{Z}_-$ belongs to $\mathcal{W}_1 \oplus \mathcal{W}_3$, then (M, g) is anti-self dual with negative scalar curvature $s = -\frac{6}{t_1}$ by Proposition 2 and [14, Theorem 4.7 (iv)]. Then conditions (27) reduce to

$$(28) \quad \mathfrak{S} \{ (-1)^n g(\mathcal{R}(U_2), Y \wedge Z - J_1 Y \wedge J_1 Z) + g(\mathcal{R}(J_2 U_2), J_1 Y \wedge Z + Y \wedge J_1 Z) \} = 0 \\ g(\mathcal{R}(J_2 V_2), J_1) = 0,$$

where the cyclic sum is over (U_2, Y, Z) , (V_2, Z, X) , (W_2, X, Y) for $J = (J_1, J_2)$, $X, Y, Z \in T_{\pi(J)}M$, $U_2, V_2, W_2 \in \mathcal{V}_{J_2}$.

For the restriction of $(H_t, \mathcal{J}^n)$ to $\mathcal{Z}_+ \times_M \mathcal{Z}_+$, $n = 1, 2$, the second identity of (28) cannot be always satisfied since $s \neq 0$. This remark proves (a). The second identity of (28) implies $\mathcal{B} = 0$ when considering the restriction of $(H_t, \mathcal{J}^n)$ to $\mathcal{Z}_+ \times_M \mathcal{Z}_-$, $n = 3, 4$. In this case, the first identity of (28) is satisfied when $\mathcal{B} = 0$. This proves (b). $\square$

Theorem 7.

- (a) $(H_t, \mathcal{J}^n)$, $n = 3$ or 4 , restricted to $\mathcal{G}_{++}$ does not belong to the class $\mathcal{W}_1$.

- (b) $(H_t, \mathcal{J}^n)$, $n = 3$ or 4 , restricted to $\mathcal{G}_{+-}$ belongs to the class $\mathcal{W}_1$ if and only if (M, g) is Einstein with positive scalar curvature s , anti-self-dual, $\mathcal{W}_-(\tau) = -\frac{s}{12}\tau$ for $\tau \in \Lambda_-^2 TM$, and $t_1 = \frac{3}{s}$.

Proof. The claim (a) directly follows from Theorem 5 (a).

By Theorems 4 (b) and 5 (b), if $(H_t, \mathcal{J}^n)$, $n = 3$ or 4 , restricted to $\mathcal{Z}_+ \times_M \mathcal{Z}_-$ belongs to $\mathcal{W}_1$, then (M, g) is Einstein, anti-self-dual, $\mathcal{R}|\Lambda_-^2 TM = 0$, and $t_1 = \frac{3}{s}$. It is easy to check by means of Proposition 1 that under these conditions

$$(D_A \Omega_{t,n})(A, B) = 0.$$

This proves (b) since the latter identity is the defining condition for the class $\mathcal{W}_1$. $\square$

Theorems 4 and 6 imply the following statement whose proof is similar to that of the preceding theorem.

Theorem 8.

- (a) $(H_t, \mathcal{J}^n)$, $n = 3$ or 4 , restricted to $\mathcal{G}_{++}$ does not belong to the class $\mathcal{W}_2$.
(b) $(H_t, \mathcal{J}^n)$, $n = 3$ or 4 , restricted to $\mathcal{G}_{+-}$ belongs to the class $\mathcal{W}_2$ if and only if (M, g) is Einstein with negative scalar curvature s , anti-self-dual, $\mathcal{W}_-(\tau) = -\frac{s}{12}\tau$ for $\tau \in \Lambda_-^2 TM$, and $t_1 = -\frac{6}{s}$.

REFERENCES

- [1] M. F. Atiyah, N. J. Hitchin, I. M. Singer, *Self-duality in four-dimensional Riemannian geometry*, Proc. Roy. Soc. London, Ser.A **362** (1978), 425-461.
- [2] A. Besse, *Einstein manifolds*, Classics in Mathematics, Springer-Verlag, 2008.
- [3] J. Davidov, O. Muškarov, *On the Riemannian curvature of a twistor space*, Acta Math. Hungarica **58**, no.3-4 (1991), 319-332.
- [4] J. Davidov, *Einstein condition and the twistor space of compatible partially complex structures*, Diff. Geom. Appl. **22** (2005), 159-179.
- [5] J. Davidov, *Harmonic almost Hermitian structures*, in S.Chiossi, A.Fino, F.Podestà, E.Musso, L.Vezzoni (Editors), *Special Metrics and Group Actions in Geometry*, Proceedings of the workshop "New perspectives in differential geometry: special metrics and quaternionic geometry", held in Rome, 16-20 November, 2015, Springer INdAM Series **23**, Springer-Verlag, 2017, pp. 129-159.
- [6] J. Davidov, *Generalized metrics and generalized twistor spaces*, Math. Z. **291** (2019), 17-46.
- [7] J. Davidov, *Product twistor spaces and Weyl geometry*, Proc. Amer. Math. Soc. **148** (2020), 3491-3506.
- [8] J. Eells, S. Salamon, *Twistorial constructions of harmonic maps of surfaces into four-manifolds*, Ann. Scuola Norm. Sup. Pisa, ser.IV, **12** (1985), 589-640.
- [9] Th. Friedrich, H. Kurke, *Compact four-dimensional self-dual Einstein manifolds with positive scalar curvature*, Math.Nachr. **106** (1982), 271-299.
- [10] P. Gauduchon, *Structures de Weyl et théorèmes d'annulation sur une variété conforme autoduale*, Ann. Scuola Norm. Sup., ser.IV, **18** (1991), 563-629.
- [11] A. Gray, L.M. Hervella, *The sixteen classes of almost Hermitian manifolds and their linear invariants*, Ann. Mat. Pure Appl. **123** (1980), 35-50.
- [12] M. Gualtieri, *Generalized complex geometry*, Ph.D. thesis, St John's College, University of Oxford, 2003, arXiv: **math.DG/0401221**.
- [13] N. Hitchin, *Generalized Calabi-Yau manifolds*, Quart. J. Math. **54** (2004), 281-308.
- [14] O. Mushkarov, *Almost Hermitian structures on twistor spaces and their type*, Atti Sem.Mat.Fis.Univ.Modena **37** (1989), 285-297.
- [15] I. M. Singer, J. A. Thorpe, *The curvature of 4-dimensional Einstein spaces*, in papers in Honor of K. Kodaira, Princeton University Press (Princeton), 1969, pp. 355-365.

- [16] J. Vilms, Totally geodesic maps, J.Diff.Geom. 4 (1970), 73-79.

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